# MODELLING AND SIMULATION OF LIQUID-VAPOR PHASE TRANSITION

A Boiling Crisis Study Contribution

Gloria Faccanoni<sup>1,2</sup> Grégoire Allaire<sup>1,2</sup> Samuel Kokh<sup>2</sup>

<sup>1</sup>École Polytechnique - CMAP

<sup>2</sup>CEA Saclay - SFME/LETR





# **BOILING CRISIS**

#### PHENOMENON

Liquid phase heated by a wall at a fixed temperature  $T^{\text{wall}}$  (pool boiling). When  $T^{\text{wall}}$  increases, we switch from a nucleate boiling to a film boiling.

Nucleate Boiling

Film Boiling





source: http://www.spaceflight.esa.int/users/fluids/TT\_boiling.htm

# OUTLINE



- 2 Numerical Method
- O Numerical Tests



# OUTLINE



- Numerical Method
- 3 Numerical Tests
- 4 Conclusion

## EULER SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = \mathbf{0}, \\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P \ \mathbb{I}) = \mathfrak{V}_{\mathrm{vf}} - \mathfrak{S}_{\mathrm{sf}}, \\ \partial_t \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \right) + \operatorname{div} \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \mathbf{u} + P \ \mathbf{u} \right) = (\mathfrak{V}_{\mathrm{vf}} - \mathfrak{S}_{\mathrm{sf}}) \cdot \mathbf{u} - \operatorname{div}(q). \end{cases}$$

- $(\mathbf{x}, t) \mapsto \rho$  specific density,
- $(\mathbf{x}, t) \mapsto \varepsilon$  specific internal energy,
- $(\mathbf{x}, t) \mapsto \mathbf{u}$  velocity;

- $(
  ho, arepsilon) \mapsto \mathfrak{V}_{\mathrm{vf}}$  volumic forces,
- $(\rho, \varepsilon) \mapsto \mathfrak{S}_{\mathrm{sf}}$  surface forces,
- $(\rho, \varepsilon) \mapsto \operatorname{div}(q)$  heat transfert.



# EULER SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = \mathbf{0}, \\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + \mathbf{P} \mathbb{I}) = \mathfrak{V}_{\mathrm{vf}} - \mathfrak{S}_{\mathrm{sf}}, \\ \partial_t \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \right) + \operatorname{div} \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \mathbf{u} + \mathbf{P} \mathbf{u} \right) = (\mathfrak{V}_{\mathrm{vf}} - \mathfrak{S}_{\mathrm{sf}}) \cdot \mathbf{u} - \operatorname{div}(q). \end{cases}$$

- $(\mathbf{x}, t) \mapsto \rho$  specific density,
- $(\mathbf{x}, t) \mapsto \varepsilon$  specific internal energy,
- $(\mathbf{x}, t) \mapsto \mathbf{u}$  velocity;

- $(\rho, \varepsilon) \mapsto \mathfrak{V}_{vf}$  volumic forces,
- $(\rho, \varepsilon) \mapsto \mathfrak{S}_{sf}$  surface forces,
- $(\rho, \varepsilon) \mapsto \operatorname{div}(q)$  heat transfert.

 $(\rho, \varepsilon) \mapsto P$  pressure law.



$$(
ho, arepsilon) \mapsto P = \left\{ egin{array}{cc} P^{ ext{liq}} & ext{if } arphi = 1; \ P^{ ext{vap}} & ext{if } arphi = 0. \end{array} 
ight.$$



$$(\rho, \varepsilon) \mapsto P = \begin{cases} P^{\text{liq}} & \text{if } \varphi = 1; \\ ??? & \text{if } 0 < \varphi < 1; \\ P^{\text{vap}} & \text{if } \varphi = 0. \end{cases}$$



$$(
ho, arepsilon) \mapsto P = \left\{ egin{array}{cc} P^{ ext{liq}} & ext{if } arphi = 1; \ P^{ ext{ff}} & ext{if } 0 < arphi < 1; \ P^{ ext{vap}} & ext{if } arphi = 0. \end{array} 
ight.$$



➡ Goal: define a global pressure law such that

- $(\rho, \varepsilon, \mathbf{u}, P)$  are continuous (3 zones)
- the interface position and the phase change are implicit (→ )
- coherence with classical thermodynamics [H. Callen]

# EOS of each Phase $\alpha = 1, 2$

$$\begin{array}{l} \tau_{\alpha} \text{ specific volume} \\ \varepsilon_{\alpha} \text{ specific internal energy} \end{array} \right\} \quad \Rightarrow \quad \mathbf{w}_{\alpha} \stackrel{\text{\tiny def}}{=} (\tau_{\alpha}, \varepsilon_{\alpha});$$

 $\mathbf{w}_{\alpha} \mapsto s_{\alpha}$  specific entropy (Hessian matrix neg. def.);

$$\begin{cases} \mathcal{T}_{\alpha} \stackrel{\text{def}}{=} \left( \frac{\partial s_{\alpha}}{\partial \varepsilon_{\alpha}} \Big|_{\tau_{\alpha}} \right)^{-1} > 0 \quad \text{temperature,} \\ P_{\alpha} \stackrel{\text{def}}{=} \mathcal{T}_{\alpha} \frac{\partial s_{\alpha}}{\partial \tau_{\alpha}} \Big|_{\varepsilon_{\alpha}} > 0 \quad \text{pressure,} \\ g_{\alpha} \stackrel{\text{def}}{=} \varepsilon_{\alpha} + P_{\alpha} \tau_{\alpha} - \mathcal{T}_{\alpha} s_{\alpha} \quad \text{free enthalpy (Gibbs potential)} \end{cases}$$

# EOS OF EACH PHASE $\alpha = 1, 2$

$$\begin{array}{c} \tau_{\alpha} \text{ specific volume} \\ \varepsilon_{\alpha} \text{ specific internal energy} \end{array} \right\} \quad \Rightarrow \quad \mathbf{w}_{\alpha} \stackrel{\text{\tiny def}}{=} (\tau_{\alpha}, \varepsilon_{\alpha});$$

 $\mathbf{w}_{\alpha} \mapsto s_{\alpha}$  specific entropy (Hessian matrix neg. def.);

# EOS OF EACH PHASE $\alpha = 1, 2$

$$\begin{array}{l} \tau_{\alpha} \text{ specific volume} \\ \varepsilon_{\alpha} \text{ specific internal energy} \end{array} \right\} \quad \Rightarrow \quad \mathbf{w}_{\alpha} \stackrel{\text{\tiny def}}{=} (\tau_{\alpha}, \varepsilon_{\alpha});$$

 $\mathbf{w}_{\alpha} \mapsto s_{\alpha}$  specific entropy (Hessian matrix neg. def.);

# EOS WITHOUT PHASE CHANGE

- $\mathbf{w} \stackrel{\text{\tiny def}}{=} y \mathbf{w}_1 + (1 y) \mathbf{w}_2;$
- y mass fraction;
- *z* volume fraction s.t.  $y\tau_1 = z\tau$ ;
- $\psi$  energy fraction s.t.  $y\varepsilon_1 = \psi\varepsilon$ .

# EOS WITHOUT PHASE CHANGE

• 
$$\mathbf{w} \stackrel{\text{\tiny def}}{=} y \mathbf{w}_1 + (1 - y) \mathbf{w}_2;$$

#### • y mass fraction;

- *z* volume fraction s.t.  $y\tau_1 = z\tau$ ;
- $\psi$  energy fraction s.t.  $y\varepsilon_1 = \psi\varepsilon$ .

#### **ENTROPY WITHOUT PHASE CHANGE**

$$\sigma \stackrel{\text{\tiny def}}{=} y s_1(\mathbf{w}_1) + (1-y) s_2(\mathbf{w}_2) = y s_1\left(\frac{z}{y}\tau, \frac{\psi}{y}\varepsilon\right) + (1-y) s_2\left(\frac{1-z}{1-y}\tau, \frac{1-\psi}{1-y}\varepsilon\right)$$

# EOS WITHOUT PHASE CHANGE

- $\mathbf{w} \stackrel{\text{\tiny def}}{=} y \mathbf{w}_1 + (1 y) \mathbf{w}_2;$
- y mass fraction;
- *z* volume fraction s.t.  $y\tau_1 = z\tau$ ;
- $\psi$  energy fraction s.t.  $y\varepsilon_1 = \psi\varepsilon$ .

#### **ENTROPY WITHOUT PHASE CHANGE**

$$\sigma \stackrel{\text{\tiny def}}{=} y s_1(\mathbf{w}_1) + (1-y) s_2(\mathbf{w}_2) = y s_1\left(\frac{z}{y}\tau, \frac{\psi}{y}\varepsilon\right) + (1-y) s_2\left(\frac{1-z}{1-y}\tau, \frac{1-\psi}{1-y}\varepsilon\right)$$







Optimality Condition:

 $P_1(z, y, \psi) = P_2(z, y, \psi)$   $P_1(z, y, \psi) = P_2(z, y, \psi)$   $g_1(z, y, \psi) = g_2(z, y, \psi)$  $z, y, \psi \in ]0, 1[^3$ 

Solution:  $(z^*, y^*, \psi^*)$ 





# DYNAMIC LIQUID-VAPOR PHASE CHANGE

#### **EULER SYSTEM**

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0, \\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + \boldsymbol{P}^{eq} \mathbb{I}) = 0 \\ \partial_t \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \right) + \operatorname{div}\left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \mathbf{u} + \boldsymbol{P}^{eq} \mathbf{u} \right) = 0 \end{cases} \text{ with } \boldsymbol{P}^{eq} \stackrel{\text{def}}{=} \frac{\boldsymbol{s}_{\varepsilon}^{eq}}{\boldsymbol{s}_{\varepsilon}^{eq}}.$$

#### **PROPERTIES** [G. ALLAIRE, G. FACCANONI, S. KOKH]

- If  $\tau_1^* \neq \tau_2^*$  and  $\varepsilon_1^* \neq \varepsilon_2^*$  (first order phase transition) then  $\mathbf{O} \ c(\mathbf{w}) > 0$ ,  $\mathbf{O} \ s_{\tau\varepsilon}^{eq}(\mathbf{w}) > 0$ 
  - Euler system: strict hyperbolicity (\u2274 p-system),
  - Riemann problem: multitude of entropic (Lax) solutions [R. Menikoff, B. J. Plohr], uniqueness of Liu solution.

# DYNAMIC LIQUID-VAPOR PHASE CHANGE

#### EULER SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0, \\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + \boldsymbol{P}^{eq} \mathbb{I}) = 0 \\ \partial_t \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \right) + \operatorname{div}\left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \mathbf{u} + \boldsymbol{P}^{eq} \mathbf{u} \right) = 0 \end{cases} \text{ with } \boldsymbol{P}^{eq} \stackrel{\text{def}}{=} \frac{\boldsymbol{s}_{\tau}^{eq}}{\boldsymbol{s}_{\varepsilon}^{eq}}.$$

#### **PROPERTIES** [G. ALLAIRE, G. FACCANONI, S. KOKH]

- If  $\tau_1^* \neq \tau_2^*$  and  $\varepsilon_1^* \neq \varepsilon_2^*$  (first order phase transition) then **1**  $c(\mathbf{w}) > 0$ ,  $\mathfrak{S}_{\tau\varepsilon}^{eq}(\mathbf{w}) > 0$ 
  - Euler system: strict hyperbolicity ( $\neq$  p-system),
  - Riemann problem: multitude of entropic (Lax) solutions [R. Menikoff,
     B. J. Plohr], uniqueness of Liu solution.

# DYNAMIC LIQUID-VAPOR PHASE CHANGE

#### EULER SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0, \\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + \mathbf{P}^{eq} \mathbb{I}) = 0 \\ \partial_t \left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \right) + \operatorname{div}\left( \rho \left( \frac{|\mathbf{u}|^2}{2} + \varepsilon \right) \mathbf{u} + \mathbf{P}^{eq} \mathbf{u} \right) = 0 \end{cases} \text{ with } \mathbf{P}^{eq} \stackrel{\text{def}}{=} \frac{\mathbf{s}_{\tau}^{eq}}{\mathbf{s}_{\varepsilon}^{eq}}.$$

#### **PROPERTIES** [G. ALLAIRE, G. FACCANONI, S. KOKH]

- If  $\tau_1^* \neq \tau_2^*$  and  $\varepsilon_1^* \neq \varepsilon_2^*$  (first order phase transition) then **1**  $c(\mathbf{w}) > 0$ , **2**  $s_{\tau\varepsilon}^{eq}(\mathbf{w}) > 0$ 
  - Euler system: strict hyperbolicity ( $\neq$  p-system),
  - Riemann problem: multitude of entropic (Lax) solutions [R. Menikoff, B. J. Plohr], uniqueness of Liu solution.

# OUTLINE



#### 2 Numerical Method

3 Numerical Tests

#### Conclusion

#### HOW TO SIMULATE THE LIU SOLUTION

- Exact Riemann Solver [A. Voß]
- Viscuous Solver (the Liu solution is the only solution that has a viscuous profile) [S. Jaouen]
- Solver(s) based on Relaxation Approach [F. Coquel, B. Perthame], [Th. Barberon, Ph. Helluy], [Ph. Helluy, N. Seguin], [F. Coquel, F. Caro, D. Jamet, S. Kokh],

 $\partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$ 

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V}) \qquad \xrightarrow{\text{Formally}} \qquad \partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$$

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V}) \qquad \xrightarrow{\text{Formally}} \qquad \partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$$

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P^{\text{eq}} \mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P^{\text{eq}})\mathbf{u}) = 0 \end{cases}$$
$$P^{\text{eq}}(\rho, \varepsilon) = \frac{s_{\varepsilon}^{\text{eq}}}{s_{\varepsilon}^{\text{eq}}}, \quad e^{\frac{\operatorname{de}}{2}} \frac{|\mathbf{u}|^2}{2} + \varepsilon$$

Formally

 $\mu \rightarrow 0$ 

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P\mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P)\mathbf{u}) = 0 \end{cases}$$

$$P(\rho,\varepsilon,z,y,\psi) = \frac{\sigma_{\tau}}{\sigma_{\varepsilon}}$$

 $\partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$ 

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$

Formally 
$$\mu \rightarrow 0$$

$$\partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$$

#### AUGMENTED SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P^{eq} \mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P^{eq})\mathbf{u}) = 0 \end{cases}$$
$$P^{eq}(\rho, \varepsilon) = \frac{s_{\varepsilon}^{eq}}{s_{\varepsilon}^{eq}}, \quad e^{\frac{\operatorname{def}}{2}} \frac{|\mathbf{u}|^2}{2} + \varepsilon$$

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$

$$\mu \rightarrow 0$$

Formall  $\mu_i \rightarrow 0$ 

$$\partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$$

#### AUGMENTED SYSTEM

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P \mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P)\mathbf{u}) = 0 \end{cases}$$

$$\underset{\mathbf{u} \in \mathcal{U}}{\underset{\mathbf{u} \in \mathcal{U}}{\underset{\mathbf{u}}{\operatorname{grad}}}} \begin{cases} \partial_t z + \mathbf{u} \cdot \operatorname{grad} z = \frac{1}{\mu_z} \left(\frac{P_2}{T_2} - \frac{P_1}{T_1}\right)\\ \partial_t y + \mathbf{u} \cdot \operatorname{grad} y = \frac{1}{\mu_y} \left(\frac{g_1}{T_1} - \frac{g_2}{T_2}\right) \frac{1}{\rho}\\ \partial_t \psi + \mathbf{u} \cdot \operatorname{grad} \psi = \frac{1}{\mu_\psi} \left(\frac{1}{T_1} - \frac{1}{T_2}\right) \varepsilon \end{cases}$$

$$P(\rho, \varepsilon, z, y, \psi) = \frac{\sigma_{\varepsilon}}{\sigma_{\varepsilon}}$$

$$\begin{array}{l} \stackrel{\text{V}}{\longrightarrow} & \begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P^{\text{eq}} \mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P^{\text{eq}})\mathbf{u}) = 0 \end{cases} \\ P^{\text{eq}}(\rho, \varepsilon) = \frac{s_t^{\text{eq}}}{s_{\varepsilon}^{\text{eq}}}, \quad e^{\frac{\operatorname{det}}{2}} \frac{|\mathbf{u}|^2}{2} + \varepsilon \end{aligned}$$

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$

$$\mu \rightarrow 0$$

Formall  $\mu_i \rightarrow 0$ 

G. Faccanoni

$$\partial_t \mathbf{U} + \operatorname{div} \mathbf{F}(\mathbf{U}) = \mathbf{0}$$

#### AUGMENTED SYSTEM

REMARK: 
$$\partial_t \psi + \mathbf{u} \cdot \mathbf{grad} \psi = 0 \rightsquigarrow T_1 = T_2$$
.

$$\begin{array}{l} \stackrel{\text{V}}{\longrightarrow} & \begin{cases} \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0\\ \partial_t(\rho \mathbf{u}) + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u} + P^{\text{eq}} \mathbb{I}) = 0\\ \partial_t(\rho e) + \operatorname{div}((\rho e + P^{\text{eq}})\mathbf{u}) = 0 \end{cases} \\ P^{\text{eq}}(\rho, \varepsilon) = \frac{s_{\varepsilon}^{\text{eq}}}{s_{\varepsilon}^{\text{eq}}}, \quad e^{\frac{\operatorname{def}}{2}} \frac{|\mathbf{u}|^2}{2} + \varepsilon \end{aligned}$$

# NUMERICAL SCHEME

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \mathbf{S}(\mathbf{V}) + \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$



Heat: 2D implicit

P, T, g equilibrium

### NUMERICAL SCHEME

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \mathbf{S}(\mathbf{V}) + \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$

**V**<sup>*n*</sup><sub>*i*</sub>

Aug. System: 5-eq. iso-T Num. Scheme: op. splitting Conv.: [G. Allaire and all.] Surf. Tens.: [J. U. Brackbill and all Heat: 2D implicit  $\mathbf{\mathfrak{G}} \ \mu_j = \mathbf{0}$   $\downarrow$   $\mathbf{R}(\mathbf{V}) = \mathbf{0}$ 

update fractions  $(y, z, \psi)$  by projecting  $\mathbf{V}_i^{n+1/2}$ onto the *P*, *T*, *g* equilibrium

### NUMERICAL SCHEME

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \mathbf{S}(\mathbf{V}) + \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$



Aug. System: 5-eq. iso-T Num. Scheme: op. splitting Conv.: [G. Allaire and all.] Surf. Tens.: [J. U. Brackbill and all.] Heat: 2D implicit update fractions  $(y, z, \psi)$  by projecting  $\mathbf{V}_i^{n+1/2}$ onto the *P*, *T*, *g* equilibrium
#### NUMERICAL SCHEME

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \mathbf{S}(\mathbf{V}) + \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$



Aug. System: 5-eq. iso-T Num. Scheme: op. splitting Conv.: [G. Allaire and all.] Surf. Tens.: [J. U. Brackbill and all.] Heat: 2D implicit update fractions  $(y, z, \psi)$  by projecting  $\mathbf{V}_i^{n+1/2}$ onto the *P*, *T*, *g* equilibrium

# NUMERICAL SCHEME

$$\partial_t \mathbf{V} + \operatorname{div} \mathbf{G}(\mathbf{V}) = \mathbf{S}(\mathbf{V}) + \frac{1}{\mu} \mathbf{R}(\mathbf{V})$$



 $(\tau, \varepsilon)$  fixed

$$( au_1,arepsilon_1, au_2,arepsilon_2,y)$$
 solution of

$$\begin{cases} g_{1}(\tau_{1}, \varepsilon_{1}) = g_{2}(\tau_{2}, \varepsilon_{2}) \\ P_{1}(\tau_{1}, \varepsilon_{1}) = P_{2}(\tau_{2}, \varepsilon_{2}) \\ T_{1}(\tau_{1}, \varepsilon_{1}) = T_{2}(\tau_{2}, \varepsilon_{2}) \\ \tau = y\tau_{1} + (1 - y)\tau_{2} \\ \varepsilon = y\varepsilon_{1} + (1 - y)\varepsilon_{2} \end{cases}$$

$$(P,T) \text{ solution of}$$

$$\begin{cases} g_1(P,T) = g_2(P,T) \\ \\ \frac{\tau - \tau_2(P,T)}{\tau_1(P,T) - \tau_2(P,T)} = \frac{\varepsilon - \varepsilon_2(P,T)}{\varepsilon_1(P,T) - \varepsilon_2(P,T)} \end{cases}$$

$$\mathsf{T}\mapsto\mathsf{P}=\mathsf{P}^{\mathrm{sat}}(\mathsf{T})pprox\mathsf{P}^{\mathrm{sat}}(\mathsf{T})$$

$$\frac{\tau - \tau_2^{\text{sat}}(T)}{\tau_1^{\text{sat}}(T) - \tau_2^{\text{sat}}(T)} = \frac{\varepsilon - \varepsilon_2^{\text{sat}}(T)}{\varepsilon_1^{\text{sat}}(T) - \varepsilon_2^{\text{sat}}(T)}$$

here 
$$\begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha}^{\operatorname{sat}}(T) \stackrel{\text{\tiny def}}{=} \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha} (P^{\operatorname{sat}})$$

#### $(\tau, \varepsilon)$ fixed

#### $( au_1,arepsilon_1, au_2,arepsilon_2,y)$ solution of

$$\begin{cases} g_1(\tau_1, \varepsilon_1) = g_2(\tau_2, \varepsilon_2) \\ P_1(\tau_1, \varepsilon_1) = P_2(\tau_2, \varepsilon_2) \\ T_1(\tau_1, \varepsilon_1) = T_2(\tau_2, \varepsilon_2) \\ \tau = y \tau_1 + (1 - y) \tau_2 \\ \varepsilon = y \varepsilon_1 + (1 - y) \varepsilon_2 \end{cases}$$

$$(P,T) \text{ solution of}$$

$$\begin{cases} g_1(P,T) = g_2(P,T) \\ \\ \frac{\tau - \tau_2(P,T)}{\tau_1(P,T) - \tau_2(P,T)} = \frac{\varepsilon - \varepsilon_2(P,T)}{\varepsilon_1(P,T) - \varepsilon_2(P,T)} \end{cases}$$

$$T \mapsto P = P^{\operatorname{sat}}(T) \approx P^{\operatorname{sat}}(T)$$

$$\frac{\tau - \tau_2^{\text{sat}}(T)}{\tau_1^{\text{sat}}(T) - \tau_2^{\text{sat}}(T)} = \frac{\varepsilon - \varepsilon_2^{\text{sat}}(T)}{\varepsilon_1^{\text{sat}}(T) - \varepsilon_2^{\text{sat}}(T)}$$

$$\operatorname{re} \left( \frac{\tau}{\varepsilon} \right)_{\alpha}^{\operatorname{sat}} (T) \stackrel{\text{\tiny def}}{=} \left( \frac{\tau}{\varepsilon} \right)_{\alpha} (P^{\operatorname{sat}}(T))$$

 $(\tau, \varepsilon)$  fixed



$$\frac{\tau - \tau_2^{\text{sat}}(T)}{\tau_1^{\text{sat}}(T) - \tau_2^{\text{sat}}(T)} = \frac{\varepsilon - \varepsilon_2^{\text{sat}}(T)}{\varepsilon_1^{\text{sat}}(T) - \varepsilon_2^{\text{sat}}(T)}$$

$$\operatorname{re} \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha}^{\operatorname{sat}}(T) \stackrel{\text{\tiny def}}{=} \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha} (P^{\operatorname{sat}}(T)$$

 $(\tau, \varepsilon)$  fixed



$$\frac{\tau - \tau_{2}^{\text{sat}}(T)}{\tau_{1}^{\text{sat}}(T) - \tau_{2}^{\text{sat}}(T)} = \frac{\varepsilon - \varepsilon_{2}^{\text{sat}}(T)}{\varepsilon_{1}^{\text{sat}}(T) - \varepsilon_{2}^{\text{sat}}(T)} \quad \text{where } \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha}^{\text{sat}}(T)^{\text{def}} \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha} (P^{\text{sat}}(T), T)$$





$$\frac{\tau - \tau_{2}^{\text{sat}}(T)}{\tau_{1}^{\text{sat}}(T) - \tau_{2}^{\text{sat}}(T)} = \frac{\varepsilon - \varepsilon_{2}^{\text{sat}}(T)}{\varepsilon_{1}^{\text{sat}}(T) - \varepsilon_{2}^{\text{sat}}(T)} \quad \text{where } \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha}^{\text{sat}}(T) \stackrel{\text{def}}{=} \begin{pmatrix} \tau \\ \varepsilon \end{pmatrix}_{\alpha} (\widehat{P}^{\text{sat}}(T), T)$$

#### TABULATED EOS

➡ Water Examples

( au, arepsilon) fixed



#### TABULATED EOS

➡ Water Examples

( au, arepsilon) fixed



approximations

#### TABULATED EOS

Water Examples

( au, arepsilon) fixed



approximations

## OUTLINE



#### 2 Numerical Method

#### O Numerical Tests





Compression of a 2D Vapor Bubble involving two Stiffened Gases for water and steam.

The piston moves towards right at constant speed  $u_p = 30$  m/s.



































Nucleation of a 2D Vapor Bubbles involving two Stiffened Gases for water and steam. The temperature of the south wall is fixed at  $T^{\text{wall}} = 310 + (600 - 310)(1 + \cos(6\pi x))/2.$ 







































































































# OUTLINE



### 2 Numerical Method

### 3 Numerical Tests



# LIQUID-VAPOR PHASE TRANSITION

### Diffuse Interface Model

- global EOS always at equilibrium (entropy maximization),
- strict hyperbolicity of the Euler system,
- uniqueness of Liu solution for the Riemann problem;
- Relaxation Approach
  - 6 (or 5) equation system with relaxation terms;

### Numerical Method

- operator splitting,
- general approximate construction of global EOS (and resolution of projection step).

# LIQUID-VAPOR PHASE TRANSITION

#### Diffuse Interface Model

- global EOS always at equilibrium (entropy maximization),
- strict hyperbolicity of the Euler system,
- uniqueness of Liu solution for the Riemann problem;

### Relaxation Approach

• 6 (or 5) equation system with relaxation terms;

### Numerical Method

- operator splitting,
- general approximate construction of global EOS (and resolution of projection step).

# LIQUID-VAPOR PHASE TRANSITION

#### Diffuse Interface Model

- global EOS always at equilibrium (entropy maximization),
- strict hyperbolicity of the Euler system,
- uniqueness of Liu solution for the Riemann problem;

### Relaxation Approach

• 6 (or 5) equation system with relaxation terms;

### Numerical Method

- operator splitting,
- general approximate construction of global EOS (and resolution of projection step).

# STIFFENED GAS FOR WATER

$$( au_{lpha},arepsilon_{lpha})\mapsto s_{lpha}=c_{v_{lpha}}\ln(arepsilon_{lpha}-q_{lpha}-\pi_{lpha} au_{lpha})+c_{v_{lpha}}(\gamma_{lpha}-1)\ln au_{lpha}+m_{lpha}$$

Phase	$c_v \left[ J/(kg \cdot K) \right]$	γ	π [Pa ]	<i>q</i> [J/kg ]	<i>m</i> [J/(kg ⋅ K) ]
Water	1816.2	2.35	10 <sup>9</sup>	$-1167.056  imes 10^{3}$	-32765.55596
Steam	1040.14	1.43	0	$2030.255  imes 10^3$	-33265.65947

TABLE: Parameters proposed by [Le Metayer] for water.

 $(P,T) \mapsto \varepsilon_{\alpha} = c_{v_{\alpha}} T \frac{P + \pi_{\alpha} \gamma_{\alpha}}{P + \pi_{\alpha}} + q_{\alpha}, \qquad (P,T) \mapsto \tau_{\alpha} = c_{v_{\alpha}} (\gamma_{\alpha} - 1) \frac{T}{P + \pi_{\alpha}}.$   $T^{i} = 278K \dots 610K, \\g_{1}(P,T^{i}) = g_{2}(P,T^{i}) \Rightarrow P^{\text{sat}}(T^{i}) \end{cases} \Rightarrow \mathfrak{A} = \left\{ \left(T^{i}, P^{\text{sat}}(T^{i})\right)\right\}_{i=0}^{83}$   $\widehat{P}^{\text{sat}} \text{ defined by using a least square approximation of } \mathfrak{A}:$   $T \mapsto P^{\text{sat}}(T) \approx \widehat{P}^{\text{sat}}(T) \stackrel{\text{def}}{=} \exp\left(\sum_{k=-8}^{k=8} a_{k} T^{k}\right)$ 

#### He Helui

### STIFFENED GAS FOR WATER

$$( au_{lpha},arepsilon_{lpha})\mapsto s_{lpha}=c_{v_{lpha}}\ln(arepsilon_{lpha}-q_{lpha}-\pi_{lpha} au_{lpha})+c_{v_{lpha}}(\gamma_{lpha}-1)\ln au_{lpha}+m_{lpha}$$

Phase	$c_v \left[ J/(kg \cdot K) \right]$	γ	$\pi$ [Pa ]	<i>q</i> [J/kg ]	<i>m</i> [J/(kg ⋅ K) ]
Water	1816.2	2.35	10 <sup>9</sup>	$-1167.056  imes 10^{3}$	-32765.55596
Steam	1040.14	1.43	0	$2030.255  imes 10^3$	-33265.65947

TABLE: Parameters proposed by [Le Metayer] for water.

 $(P,T) \mapsto \varepsilon_{\alpha} = c_{v_{\alpha}} T \frac{P + \pi_{\alpha} \gamma_{\alpha}}{P + \pi_{\alpha}} + q_{\alpha}, \qquad (P,T) \mapsto \tau_{\alpha} = c_{v_{\alpha}} (\gamma_{\alpha} - 1) \frac{T}{P + \pi_{\alpha}}.$   $T^{i} = 278K \dots 610K, \\g_{1}(P,T^{i}) = g_{2}(P,T^{i}) \Rightarrow P^{\text{sat}}(T^{i}) \end{cases} \Rightarrow \mathfrak{A} = \left\{ \left(T^{i}, P^{\text{sat}}(T^{i})\right) \right\}_{i=0}^{83}$   $\widehat{P}^{\text{sat}} \text{ defined by using a least square approximation of } \mathfrak{A}:$   $T \mapsto P^{\text{sat}}(T) \approx \widehat{P}^{\text{sat}}(T) \stackrel{\text{def}}{=} \exp\left( \left( \sum_{k=0}^{k=8} - \alpha T^{k} \right) \right)$ 

# STIFFENED GAS FOR WATER

$$( au_{lpha},arepsilon_{lpha})\mapsto s_{lpha}=c_{v_{lpha}}\ln(arepsilon_{lpha}-q_{lpha}-\pi_{lpha} au_{lpha})+c_{v_{lpha}}(\gamma_{lpha}-1)\ln au_{lpha}+m_{lpha}$$

Phase	$c_v \left[ J/(kg \cdot K) \right]$	γ	$\pi$ [Pa ]	<i>q</i> [J/kg ]	<i>m</i> [J/(kg ⋅ K) ]
Water	1816.2	2.35	10 <sup>9</sup>	$-1167.056  imes 10^{3}$	-32765.55596
Steam	1040.14	1.43	0	$2030.255  imes 10^3$	-33265.65947

TABLE: Parameters proposed by [Le Metayer] for water.

 $(P,T) \mapsto \varepsilon_{\alpha} = c_{\nu_{\alpha}} T \frac{P + \pi_{\alpha} \gamma_{\alpha}}{P + \pi_{\alpha}} + q_{\alpha}, \qquad (P,T) \mapsto \tau_{\alpha} = c_{\nu_{\alpha}} (\gamma_{\alpha} - 1) \frac{T}{P + \pi_{\alpha}}.$   $T^{i} = 278K \dots 610K, \\g_{1}(P,T^{i}) = g_{2}(P,T^{i}) \Rightarrow P^{\text{sat}}(T^{i}) \end{cases} \Rightarrow \mathfrak{A} = \left\{ (T^{i}, P^{\text{sat}}(T^{i})) \right\}_{i=0}^{83}$   $\widehat{P}^{\text{sat}} \text{ defined by using a least square approximation of } \mathfrak{A}:$   $T \mapsto P^{\text{sat}}(T) \approx \widehat{P}^{\text{sat}}(T) \stackrel{\text{def}}{=} \exp\left(\sum_{k=0}^{k=0} \alpha a_{k} T^{k}\right)$ 

# STIFFENED GAS FOR WATER

$$( au_{lpha},arepsilon_{lpha})\mapsto s_{lpha}=c_{v_{lpha}}\ln(arepsilon_{lpha}-q_{lpha}-\pi_{lpha} au_{lpha})+c_{v_{lpha}}(\gamma_{lpha}-1)\ln au_{lpha}+m_{lpha}$$

Phase	$c_v \left[ J/(kg \cdot K) \right]$	γ	π [Pa ]	<i>q</i> [J/kg ]	<i>m</i> [J/(kg ⋅ K) ]
Water	1816.2	2.35	10 <sup>9</sup>	$-1167.056  imes 10^{3}$	-32765.55596
Steam	1040.14	1.43	0	$2030.255  imes 10^3$	-33265.65947

TABLE: Parameters proposed by [Le Metayer] for water.

 $\begin{aligned} (P,T) \mapsto \varepsilon_{\alpha} &= c_{v_{\alpha}} T \frac{P + \pi_{\alpha} \gamma_{\alpha}}{P + \pi_{\alpha}} + q_{\alpha}, \qquad (P,T) \mapsto \tau_{\alpha} &= c_{v_{\alpha}} (\gamma_{\alpha} - 1) \frac{T}{P + \pi_{\alpha}}. \\ T^{i} &= 278K \dots 610K, \\ g_{1}(P,T^{i}) &= g_{2}(P,T^{i}) \Rightarrow P^{\text{sat}}(T^{i}) \end{cases} \\ \geqslant \mathfrak{A} &= \left\{ \left( T^{i}, P^{\text{sat}}(T^{i}) \right) \right\}_{i=0}^{83} \\ \widehat{P}^{\text{sat}} \text{ defined by using a least square approximation of } \mathfrak{A}: \end{aligned}$ 

$$T \mapsto P^{\mathrm{sat}}(T) \approx \widehat{P}^{\mathrm{sat}}(T) \stackrel{\mathrm{def}}{=} \exp\left(\sum_{k=-8}^{k=8} a_k T^k\right)$$

$$T^{i} = 278K \dots 610K,$$

$$\varepsilon_{\alpha}^{\text{sat}}(T^{i}), \ \tau_{\alpha}^{\text{sat}}(T^{i}) \text{ found in the tables } \} \Rightarrow \begin{cases} \mathfrak{A} = \left\{ \left(T_{i}, \frac{1}{\varepsilon_{\text{vap}}^{\text{sat}}(T_{i})}\right) \right\}_{i} \\ \mathfrak{B} = \left\{ \left(T_{i}, \frac{\varepsilon_{\text{int}}^{\text{sat}}(T_{i})}{\varepsilon_{\text{vap}}^{\text{sat}}(T_{i})}\right) \right\}_{i} \\ \mathfrak{C} = \left\{ \left(T_{i}, \frac{1}{\varepsilon_{\text{vap}}^{\text{sat}}(T_{i})}\right) \right\}_{i} \\ \mathfrak{D} = \left\{ \left(T_{i}, \frac{\tau_{\text{sat}}^{\text{sat}}(T_{i})}{\varepsilon_{\text{vap}}^{\text{sat}}(T_{i})}\right) \right\}_{i} \end{cases}$$

 $\hat{\epsilon}_{\alpha}^{sat}$  and  $\hat{\tau}_{\alpha}^{sat}$  defined by using a least square approximation of  $\mathfrak{A}, \mathfrak{B}, \mathfrak{C}$  and  $\mathfrak{D}$ :

$$T \mapsto \varepsilon_{\mathrm{vap}}^{\mathrm{sat}} \approx \widehat{\varepsilon}_{\mathrm{vap}}^{\mathrm{sat}} \stackrel{\text{def}}{=} \frac{1}{\sum_{k=0}^{6} a_{k} T^{k}} \qquad T \mapsto \varepsilon_{\mathrm{liq}}^{\mathrm{sat}} \approx \widehat{\varepsilon}_{\mathrm{liq}}^{\mathrm{sat}} \stackrel{\text{def}}{=} \widehat{\varepsilon}_{\mathrm{vap}}^{\mathrm{sat}}(T) \sum_{k=0}^{6} b_{k} T^{k}$$
$$T \mapsto \tau_{\mathrm{vap}}^{\mathrm{sat}} \approx \widehat{\tau}_{\mathrm{vap}}^{\mathrm{sat}} \stackrel{\text{def}}{=} \frac{1}{\sum_{k=0}^{8} c_{k} T^{k}} \qquad T \mapsto \tau_{\mathrm{liq}}^{\mathrm{sat}} \approx \widehat{\tau}_{\mathrm{liq}}^{\mathrm{sat}} \stackrel{\text{def}}{=} \widehat{\tau}_{\mathrm{vap}}^{\mathrm{sat}}(T) \sum_{k=0}^{9} d_{k} T^{k}$$